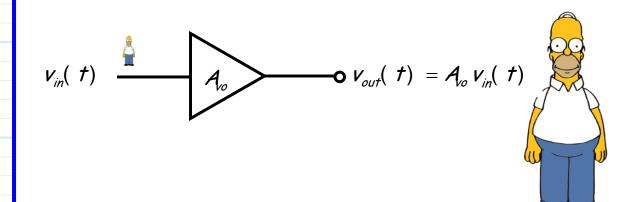
<u>Amplifiers</u>

An ideal amplifier takes an input signal and reproduces it exactly at its output, only with a larger magnitude!



where A_{vo} is the open-circuit voltage gain of the amplifier.

Now, let's express this result using our knowledge of linear circuit theory !

Recall, the output $v_{out}(t)$ of a linear device can be determined by **convolving** its input $v_{in}(t)$ with the device **impulse response** g(t):

$$v_{out}(t) = \int g(t-t')v_{in}(t')dt'$$

The impulse response for the **ideal** amplifier would therefore be:

$$g(t) = A_{vo} \delta(t)$$

so that:

$$V_{out}(t) = \int_{-\infty}^{t} g(t-t') V_{in}(t') dt'$$
$$= \int_{-\infty}^{t} A_{o} \delta(t-t') V_{in}(t') dt'$$
$$= A_{o} V_{in}(t)$$

We can alternatively represent the ideal amplifier response in the **frequency domain**, by taking the **Fourier Transform** of the impulse response:

$$T(\omega) = \int_{-\infty}^{\infty} g(t)e^{-j\omega t} dt$$
$$= \int_{-\infty}^{\infty} A_{\omega} \delta(t)e^{-j\omega t} dt$$
$$= A_{\omega} + j0$$

This result, although simple, has an interesting interpretation. It means that the amplifier exhibits gain of A_{vo} for sinusoidal signals of **any** and **all** frequencies!

$$|\mathcal{T}(\omega)| \uparrow$$

Moreover, the ideal amplifier does not alter the **relative phase** of the sinusoidal signal (i.e., no phase shift).

In other words, if:

$$V_{in}(t) = \cos(\omega t)$$

then at the output of the ideal amplifier we shall see:

$$\mathbf{v}_{out}(t) = |\mathcal{T}(\omega)|\cos(\omega t + \angle \mathcal{T}(\omega))$$
$$= \mathcal{A}_{o}\cos(\omega t)$$

BUT, there is one **big** problem with an ideal amplifier:

They are impossible to build !!

- Q: Why is that ??
- A: Two reasons:
 - a) An ideal amplifier has **infinite** bandwidth.
 - b) An ideal amplifier has **zero** delay.

Not gonna happen !

Let's look at this **second** problem first. The ideal amplifier impulse response $g(t) = A_{\omega} \delta(t)$ means that the signal at the output occurs **instantaneously** with the signal at the input.

This of course **cannot** happen, as it takes some small, but nonzero amount of **time** for the signal to propagate through the amplifier. A more **realizable** amplifier impulse response is:

$$g(t) = \mathcal{A}_{vo} \, \delta(t-\tau)$$

resulting in an amplifier output of:

$$\begin{aligned}
\mathcal{V}_{out}(t) &= \int_{-\infty} \mathcal{g}(t-t') \, \mathcal{V}_{in}(t') dt' \\
&= \int_{-\infty}^{t} \mathcal{A}_{o} \, \delta(t-\tau-t') \, \mathcal{V}_{in}(t') dt' \\
&= \mathcal{A}_{vo} \, \mathcal{V}_{in}(t-\tau)
\end{aligned}$$

In other words, the output is both an amplified and **delayed** version of the input.

* Ideally, this delay does not **distort** the signal, as the output will have the same form as the input.

* Moreover, the delay for electronic devices such as amplifiers is very small in comparison to human time scales (i.e., $\tau \ll 1$ second).

* Therefore, propagation delay τ is generally **not** considered a **problem** for most amplifier applications.

Let's examine what this delay means in the frequency domain.

Evaluating the Fourier Transform of this **modified** impulse response gives:

$$T(\omega) = \int_{-\infty}^{\infty} g(t) e^{-j\omega t} dt$$
$$= \int_{-\infty}^{\infty} A_{\omega} \delta(t - \tau) e^{-j\omega t} dt$$
$$= A_{\omega} \cos(\omega \tau) + j A_{\omega} \sin(\omega \tau)$$
$$= A_{\omega} e^{j\omega \tau}$$

We see that, as with the ideal amplifier, the magnitude $|T(\omega)| = A_{\omega}$. However, the relative **phase** is now a linear function of frequency:

$$\angle T(\omega) = \omega \tau$$

As a result, if $v_{in}(t) = \cos(\omega t)$, the output signal will be:

$$\mathbf{v}_{out}(t) = |T(\omega)| \cos(\omega t - \angle T(\omega))$$
$$= \mathcal{A}_{c} \cos(\omega t - \omega \tau)$$

In other words, the output signal of a **real** amplifier is **phase shifted** with respect to the input.

6/7

In general, the amplifier phase shift $\angle T(\omega)$ will not be a perfectly linear function (i.e., $\angle T(\omega) \neq \omega \tau$), but instead will be a more general function of frequency ω .

However, if the phase function $\angle T(\omega)$ becomes too "nonlinear", we find that signal **dispersion** can result—the output signal can be **distorted**!

Now, let's examine the **first problem** with the ideal amplifier. This problem is best discussed in the **frequency** domain.

We discovered that the **ideal** amplifier has a frequency response of $|\mathcal{T}(\omega)| = \mathcal{A}_{\omega}$. Note this means that the amplifier gain is \mathcal{A}_{vo} for all frequencies $0 < \omega < \infty$ (D.C. to daylight !).

The bandwidth of the ideal amplifier is therefore infinite !

* Since every electronic device will exhibit **some** amount of inductance, capacitance, and resistance, every device will have a **finite** bandwidth.

* In other words, there will be frequencies ω where the device does **not work**!

* From the standpoint of an amplifier, "not working" means $|T(\omega)| \ll A_o$ (i.e., **low gain**).

Amplifiers will therefore have **finite** bandwidths.

*

There is a range of frequencies ω between $\omega_{\!\scriptscriptstyle L}$ and $\omega_{\!\scriptscriptstyle H}$ where

7/7